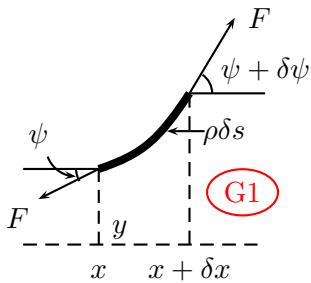


Answers and Mathscheme

1.



$$N2 \Rightarrow F\delta\psi \approx \rho\delta s \frac{\partial^2 y}{\partial t^2} \quad \text{M1A1}$$

But oscillations small $\Rightarrow \delta s \approx \delta x$, $\psi \approx \frac{\partial y}{\partial x}$

$$\Rightarrow \delta\psi \approx \frac{\partial^2 y}{\partial x^2} \delta x \Rightarrow \quad y_{xx} = \frac{1}{c^2} y_{tt}$$

M1A1

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- Write $u = x - ct$, $v = x + ct$ M1

$$y_x = y_u + y_v \Rightarrow y_{xx} = \left(\frac{\partial}{\partial u} + \frac{\partial}{\partial v} \right) (y_u + y_v) = y_{uu} + 2y_{uv} + y_{vv}$$

M2A1+A1

$$y_t = -cy_u + cy_v \Rightarrow y_{tt} = c^2 \left(\frac{\partial}{\partial u} - \frac{\partial}{\partial v} \right) (y_u - y_v) = c^2 (y_{uu} - 2y_{uv} + y_{vv})$$

- Substitute in PDE $\Rightarrow y_{uv} = 0$ M1A1 VERIFICATION M1A2 ONLY

$$y_{uv} = 0 \Rightarrow y_u = f_*(u) \quad \text{M1A1} \quad \Rightarrow y = \int^u f_*(s) ds + g(v)$$

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$$\Rightarrow \quad y(x, t) = f(x - ct) + g(x + ct)$$

M1A1

- $y = f(x - ct) + g(x + ct)$

$$t = 0 : f(x) + g(x) = 0 \quad \text{M1A1} \quad \text{and} \quad -cf'(x) + cg'(x) = V \cos 2kx$$

$$\Rightarrow -f(x) + g(x) = \frac{V}{2kc} \sin 2kx \quad (+d) \rightarrow \text{no effect on } y \quad \text{M1A1}$$

$$\therefore f(x) = -\frac{V}{4kc} \sin 2kx, g(x) = \frac{V}{4kc} \sin 2kx \quad \text{A1} \leftarrow \text{FOR EITHER}$$

$$\therefore y(x, t) = \frac{V}{4kc} \{ \sin[2k(x + ct)] - \sin[2k(x - ct)] \}$$

M1A1 FT to end bits

$$y(x, t) = \frac{V}{2kc} \cos(2kx) \sin(2kct) \Rightarrow y(x, t) = 0 \text{ when } 2kx = (n + \frac{1}{2})\pi$$

$$\Rightarrow \quad x = \frac{(2n + 1)\pi}{4k} \quad \text{A1}$$

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Standard waves with wavelength $\frac{\pi}{k}$ and period $\frac{\pi}{kc}$

A1

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2. Let $y_n = \sin\left(\frac{n\pi x}{l}\right) \cos\left(\frac{n\pi c t}{l}\right)$

M1A1

M1A1

Then $c^2 y_n'' - \ddot{y}_n = -\frac{n^2 \pi^2 c^2}{l^2} y_n - \left(-\frac{n^2 \pi^2 c^2}{l^2}\right) y_n = 0$ so linearity \Rightarrow PDE satisfied

(a) ok since $y_n(0, t) = y_n(l, t) = 0$ M1A1

(c) ok since $\dot{y}_n|_{t=0} = 0$ because $\sin\left(\frac{n\pi x}{l}\right) = 0$ at $t = 0$ M1A1

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We need $\frac{h}{l^2} x(l-x) = \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{l} \Rightarrow \frac{a_n l}{2} = \int_0^l \frac{h}{l^2} x(l-x) \sin \frac{n\pi x}{l} dx$

M1A1

M1A1

$$\begin{aligned} \Rightarrow a_n &= \frac{2h}{l^3} \left[\left[-\frac{l}{n\pi} \cos\left(\frac{n\pi x}{l}\right) x(l-x) \right]_0^l + \frac{l}{n\pi} \int_0^l (l-2x) \cos \frac{n\pi x}{l} dx \right] \\ &= \frac{2h}{n\pi l^2} \left[\left[\frac{l}{n\pi} \sin\left(\frac{n\pi x}{l}\right) (l-2x) \right]_0^l + \frac{2l}{n\pi} \int_0^l \sin \frac{n\pi x}{l} dx \right] \\ &= \frac{4h}{(n\pi)^3} \left[\cos\left(\frac{n\pi x}{l}\right) \right]_0^l = \begin{cases} 8h/(n\pi)^3 & (\text{n odd}) \\ 0 & (\text{n even}) \end{cases} \end{aligned}$$

M1A1 + A1

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$V = \frac{1}{2} \rho c^2 \int_0^l \left(\frac{\partial y}{\partial x} \right)^2 dx = \frac{1}{2} \rho c^2 \frac{\pi^2}{l^2} \sum_{n=0}^{\infty} \cos^2 \frac{n\pi c t}{l} n^2 \int_0^l a_n^2 \cos^2 \frac{n\pi x}{l} dx$

M1A1 + A1

M1A1 if M0 here, no more maths

by orthogonality

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$$\Rightarrow V = \frac{\rho \pi^2 c^2}{4l} \sum_{n=0}^{\infty} a_n^2 n^2 \cos^2 \frac{n\pi c t}{l} = \frac{16\rho h^2 c^2}{\pi^4 l} \sum_{m=0}^{\infty} \frac{\cos^2 \left\{ \frac{(2m+1)\pi c t}{l} \right\}}{(2m+1)4}$$

similar examples seen.

M1A1

M1

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3.

$$(i) \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} \left((\rho_0 + [\rho - \rho_0])u \right) + \frac{\partial}{\partial y} \left((\rho_0 + [\rho - \rho_0])v \right) + \frac{\partial}{\partial z} \left((\rho_0 + [\rho - \rho_0])w \right) = 0 \quad M1/A1$$

$$\frac{\partial \rho}{\partial t} + \rho_0 \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) + \frac{\partial}{\partial x} \underbrace{[(\rho - \rho_0)u + \dots]}_{M1A1} = 0$$

2nd order in smallness

$$\Rightarrow \rho_t + \rho_0(u_x + v_y + w_z) \approx 0 \quad A1$$

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$$(ii) \rho_{tt} = -\rho_0 \frac{\partial}{\partial x}(u_t) - \rho_0 \frac{\partial}{\partial y}(v_t) - \rho_0 \frac{\partial}{\partial z}(w_t) \quad A1/M1$$

N2 for each component M1/A1

$$\rho_0 u_t = -p_x, \quad \rho_0 v_t = -p_y, \quad \rho_0 w_t = -p_z$$

$$\Rightarrow \rho_{tt} = p_{xx} + p_{yy} + p_{zz} \quad A1$$

$$\text{But } p = f(\rho) \Rightarrow (p - p_0) = \frac{df}{d\rho}|_{\rho_0}(\rho - \rho_0) = c^2(\rho - \rho_0) \quad M2/A1$$

$$\Rightarrow p_{xx} = c^2 \rho_{xx}, \quad A1 \quad p_{yy} = c^2 \rho_{yy}, \quad p_{zz} = c^2 \rho_{zz}$$

$$\Rightarrow \rho_{tt} = c^2(\rho_{xx} + \rho_{yy} + \rho_{zz}) \quad A1$$

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$$(iii) \rho = \rho_0(1 + s), \quad u = \phi_x, \quad v = \phi_y, \quad w = \phi_z$$

$$\Rightarrow \rho_0 s_t = -\rho_0(\phi_{xx} + \phi_{yy} + \phi_{zz})$$

$$s_t = -(\phi_{xx} + \phi_{yy} + \phi_{zz}) \quad (1) \quad M1/A1$$

From N2

$$\frac{\partial}{\partial x}(\phi_t) = -c^2 s_x \Rightarrow \frac{\partial}{\partial x}(c^2 s + \phi_t) = 0 \Rightarrow c^2 s + \phi_t = F(t) \quad A1$$

$$\text{As } |x| \rightarrow \infty, \text{ } \# \text{ disturbance} \Rightarrow F(t) = 0 \quad M1/A1$$

$$\Rightarrow c^2 s = -\phi_t \quad (2) \quad A1$$

From (1) & (2)

$$c^2 s_t = -\phi_{tt} = -c^2(\phi_{xx} + \phi_{yy} + \phi_{zz}) \quad M1$$

$$\Rightarrow \phi_{tt} = c^2(\phi_{xx} + \phi_{yy} + \phi_{zz}) \quad A1$$

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4. (i)

(a): vertical component of velocity approaches zero at great depth. M2

(b): vertical component of velocity of free surface is consistent with
that prescribed by η . M2

(c): pressure is continuous at free surface. M2

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(ii) Try

$\phi = f(z) \sin(kx - \omega t)$ to be consistent with (ii) M1/A1

$$f'' - k^2 f = 0 \quad \text{A1}$$

$$f = A e^{kz} + B e^{-kz} \quad \text{A1/M1}$$

$$(a) \Rightarrow B = 0 \quad \text{A1}, \quad \phi = A e^{kz} \sin(kx - \omega t) \quad \text{A1}$$

$$(b) \Rightarrow A k \sin(kx - \omega t) = a \omega \sin(kx - \omega t) \Rightarrow A = \frac{a \omega}{k} \quad \text{A1}$$

$$\phi = \frac{a \omega}{k} e^{kz} \sin(kx - \omega t) \quad \text{A1}$$

$$(c) - \frac{a \omega^2}{k} \cos(kx - \omega t) + g a \cos(kx - \omega t) = 0 \quad \text{A1/M1}$$

$$\Rightarrow \omega^2 = gk \quad \text{A1}$$

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$$c = \frac{\omega}{k} = \sqrt{\frac{g}{q}} \quad \text{M1/A1}$$

$$c_g = \frac{d\omega}{dk} = \frac{d}{dk} \sqrt{\frac{g}{k}} = \frac{1}{2} \sqrt{\frac{g}{k}} \quad \text{M1/A1} \Rightarrow c_g = \frac{1}{2} c. \quad \text{A1}$$

Energy is always propagating with c_g . M1

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5. (i) The associated equations are $\frac{dx}{y} = \frac{dy}{x} = \frac{dz}{xy}$ M1A1

$$\therefore xdx = ydy \Rightarrow x^2 - y^2 = \alpha \quad \text{M1A1}$$

$$\text{Then } \frac{dz}{dx} = x \Rightarrow z = \frac{1}{2}x^2 + \beta \quad \text{span style="border: 1px solid red; border-radius: 50%; padding: 2px;">M1A1$$

Hence $z - \frac{1}{2}x^2$ is a function of $(x^2 - y^2)$ \Rightarrow GS is $z = \frac{1}{2}x^2 + F(x^2 - y^2)$ M1A1

For $x^2 \geq y^2$, $z = e^{-x^2}$, $y = 0$

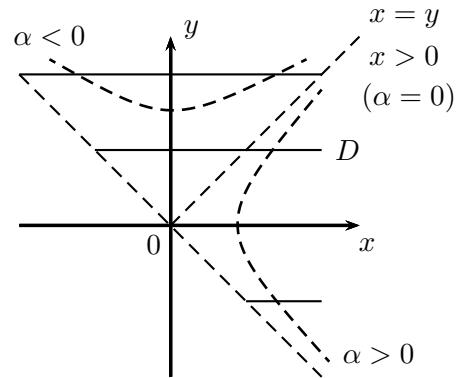
$$\Rightarrow F(x^2) = e^{-x^2} - \frac{1}{2}x^2 \quad \text{span style="border: 1px solid red; border-radius: 50%; padding: 2px;">M2A1$$

$$\Rightarrow z = \frac{1}{2}y^2 + e^{-(x^2-y^2)} \quad (x^2 \geq y^2)$$

For $x^2 \leq y^2$, $z = e^{-y^2}$, $x = 0$ needed for A1

$$\Rightarrow F(-y^2) = e^{-y^2} \quad \text{span style="border: 1px solid red; border-radius: 50%; padding: 2px;">M1A1} \quad \text{In either order}$$

$$\Rightarrow z = \frac{1}{2}x^2 + e^{(x^2-y^2)} \quad (x^2 \leq y^2) \quad \text{span style="border: 1px solid red; border-radius: 50%; padding: 2px;">B1$$



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(ii) We note that either approach gives $z = \frac{1}{2}x^2 + 1$ on $x = y, x > 0$ (*)

The characteristics that cross $x = 0, y \geq 0$ and $y = 0, x \geq 0$ cannot be extended across $x + y = 0$, so D is $x + y \geq 0$ M1A1

(*) shows z is continuous everywhere in D . M1

B1

The only curve on which discontinuities in z_x, z_y can exist is $y = x (x \geq 0)$

By calculation we find that as $x \rightarrow y$ in $x^2 > y^2$, $z_x \rightarrow -2x, z_y \rightarrow 3x$, and that as $x \rightarrow y$ in $x^2 < y^2$, $z_x \rightarrow 3x, z_y \rightarrow -2x$.

7 \Rightarrow jumps in z_x, z_y of magnitudes $\pm 5x$

M1A1

(iii) Next:

$$\text{In } + \delta z^+ = \frac{\partial z^+}{\partial x} \delta x + \frac{\partial z^+}{\partial y} \delta y, P \frac{\partial z^+}{\partial x} + Q \frac{\partial z^+}{\partial y} = R$$

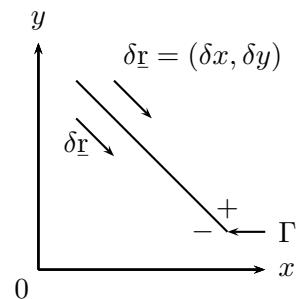
$$\text{In } - \delta z^- = \frac{\partial z^-}{\partial x} \delta x + \frac{\partial z^-}{\partial y} \delta y, P \frac{\partial z^-}{\partial x} + Q \frac{\partial z^-}{\partial y} = R$$

Since z is cont., $\delta z^+ = \delta z^-$ M1A1 + A1

$$\therefore \delta x \left[\frac{\partial z}{\partial x} \right]_+^+ + \delta y \left[\frac{\partial z}{\partial y} \right]_+^+ = 0 \text{ and } P \left[\frac{\partial z}{\partial x} \right]_+^+ + Q \left[\frac{\partial z}{\partial y} \right]_+^+ = 0$$

The necessary condition for $\left[\frac{\partial z}{\partial x} \right]_+^+$ and $\left[\frac{\partial z}{\partial y} \right]_+^+$ to be non-zero is

$$\frac{\delta x}{\delta y} = \frac{P}{Q} \Rightarrow \left(\frac{dy}{dx} = \frac{Q}{P} \right) \quad \text{span style="border: 1px solid red; border-radius: 50%; padding: 2px;">M1A1$$



BW

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